

i) Conversion of  $L_{\text{orbit}} \rightarrow L_{\text{rotation}}$

1) Before collapse all  $L$  in orbit,  $L_{\text{orb}}$

After collapse  $L$  in orbit + in rotation of planet:  $L'_{\text{orb}} + L_{\text{rot}}$

$$\therefore L_{\text{orb}} = L'_{\text{orb}} + L_{\text{rot}}$$

$$L_{\text{rot}} = L_{\text{orb}} - L'_{\text{orb}} \quad (1)$$

$$2) L'_{\text{orb}} = m_p \Omega_k a^2 \quad (2)$$

$$3) L_{\text{rot}} = \frac{2}{5} m_p R_p^2 \omega \quad (3)$$

angular frequency

assuming rigid body rotation

assuming uniform density

4) Integrate to find orbital angular momentum contained in collapsing patch of disk,  $L_{\text{orb}}$

$$dL_{\text{orb}} = dm \Omega_k(r) r^2$$

$$dm = \Sigma dA$$

$$dA = 2\pi r dr$$

$$\therefore dm = \Sigma 2\pi r dr$$

$$\text{And } \Omega_n = \left( \frac{GM_+}{r^3} \right)^{1/2} = \mu^{1/2} r^{-3/2}$$

$$\text{So, } dL_{ob} = \mu^{1/2} r^{-3/2} \cdot r^2 \cdot \Sigma \cdot 2\pi r dr$$

$$= \mu^{1/2} 2\pi \Sigma r^{3/2} dr$$

$$L_{ob} = \mu^{1/2} 2\pi \Sigma \int_{a-l_e}^{a+l_e} r^{3/2} dr$$

$$= \frac{2}{5} \mu^{1/2} 2\pi \Sigma \left[ (a+l_e)^{5/2} - (a-l_e)^{5/2} \right] \quad (2)$$

Sub (2), (3), (4)  $\rightarrow$  (1)

$$\frac{\Sigma}{5} M_p l_p^2 \omega = \frac{4}{5} \mu^{1/2} M_p \left( \frac{(a+l_e)^{5/2} - (a-l_e)^{5/2}}{4al_e} \right) - M_p \mu^{1/2} a^{1/2}$$

$$l_p^2 \omega = 2\mu^{1/2} \left[ \frac{a^{3/2}}{4l_e} \left[ \left(1 + \frac{l_e}{a}\right)^{5/2} - \left(1 - \frac{l_e}{a}\right)^{5/2} \right] \right] - \mu^{1/2} a^{1/2}$$

using:  $\Sigma = \frac{M_p}{\pi} \frac{1}{(a+l_e)^2 - (a-l_e)^2} = \frac{M_p}{\pi} \frac{1}{4al_e}$

$$\text{Now, } \left(1 + \frac{l_e}{a}\right)^{5/2} \rightarrow 1 + \frac{5}{2} \frac{l_e}{a} \dots$$

$$\left(1 - \frac{l_e}{a}\right)^{5/2} \rightarrow 1 - \frac{5}{2} \frac{l_e}{a} \dots$$

$$\frac{f^{(n)}(a) (x-a)^n}{n!}$$

$$\therefore R_p \omega \sim \frac{\mu^{1/2} a^{3/2}}{2 R_p} \left( 1 + \frac{5}{2} \frac{R_p}{a} - 1 + \frac{5}{2} \frac{R_p}{a} \right) \sim \mu^{1/2} a^{1/2}$$

$$\sim \frac{5}{2} \mu^{1/2} a^{1/2} \sim \mu a^{1/2} \quad (5)$$

$$\sim \frac{3}{2} \mu^{1/2} a^{1/2}$$

$$\text{So } \omega \sim \frac{3}{2} \frac{\mu^{1/2} a^{1/2}}{R_p^2}$$

only depend on  $M_*$ , orbital radius  
& final radius of planet

$$\text{ii) } P_{\text{day}} = \omega^{-1} \sim \left[ \frac{3}{2} \frac{(2\pi \times 10^{-11} \times 10^{30} \frac{\pi}{2} \times 10^{11})^{1/2}}{(2\pi \times 10^7)^2} \right]^{-1}$$

$$\sim \left[ \frac{10^{15}}{10^{15}} \right]^{-1}$$

$$\sim 1 \text{ s}$$

iii) Compare with self gravity

$$F_c = \frac{mv^2}{r} = \frac{m\omega^2 r^2}{r} \quad \longleftrightarrow \quad F_g = \frac{GM_p m}{r^2}$$
$$= m\omega^2 r$$

$$\frac{F_g}{F_c} = \frac{GM_p m}{r^2} \cdot \frac{1}{m\omega^2 r}$$

$$= \frac{GM_p}{\omega^2 r^3} \sim \frac{2\pi \times 10^{-11} \cdot 10^{-2} \cdot 10^{36}}{1 \cdot (2\pi \cdot 10^9)^3}$$

$$\sim \frac{10^{23}}{10^{27}}$$

$$\sim 10^{-5}$$

$\therefore$  planet would tear apart under its own rotation.

iv) Return to (5)

$$R_p^2 \omega \sim \underbrace{\frac{5}{2} \mu^{\frac{1}{2}} a^{\frac{1}{2}} - \mu^{\frac{1}{2}} a^{\frac{1}{2}}}_{L \text{ disk}}$$

$$\omega \sim \frac{\mu^{\frac{1}{2}} a^{\frac{1}{2}} \left( \frac{5}{2} f - 1 \right)}{R_p^2} \quad \parallel \text{introduce } L \text{ efficiency factor}$$

$$\text{Now, } \frac{F_g}{F_c} \geq 1 \quad \therefore \frac{GM_p m}{R_p^2} \cdot \frac{1}{m \omega_{\max}^2 R_p} \geq 1$$

$$\frac{GM_p}{R_p^3} \cdot \frac{R_p^4}{GM_* a \left( \frac{5}{2} f_{\max} - 1 \right)} \geq 1$$

$$\text{So } \frac{5}{2} f_{\max} = \frac{M_p}{M_*} \cdot \frac{R_p}{a} + 1$$

$$f_{\max} = \frac{2}{5} \left( \frac{M_p}{M_*} \cdot \frac{R_p}{a} + 1 \right)$$

$$\sim \frac{2}{5} \left( 10^{-3} \cdot \frac{10^7}{10^{11}} + 1 \right)$$

$$\sim \frac{2}{5} //$$

∴ at most 40% of angular momentum of disk can be kept in planet  
and this is almost entirely in orbit, not in planet's rotation: //

⇒ jets, winds, viscous spreading of planet's sub-disk, and  
magnetic torques on disk from planet must combine to  
allow accretion: //